

Механизмы переноса заряда в  
неупорядоченных твердых  
матрицах, изготовленных на  
основе органических материалов

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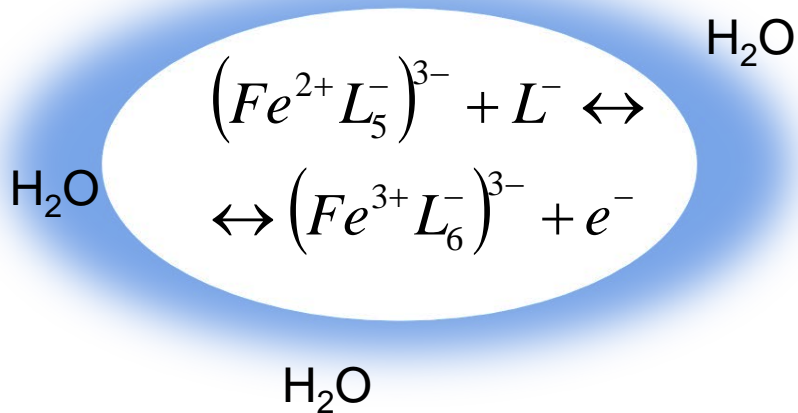
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# Ordinary ET (polar liquid)



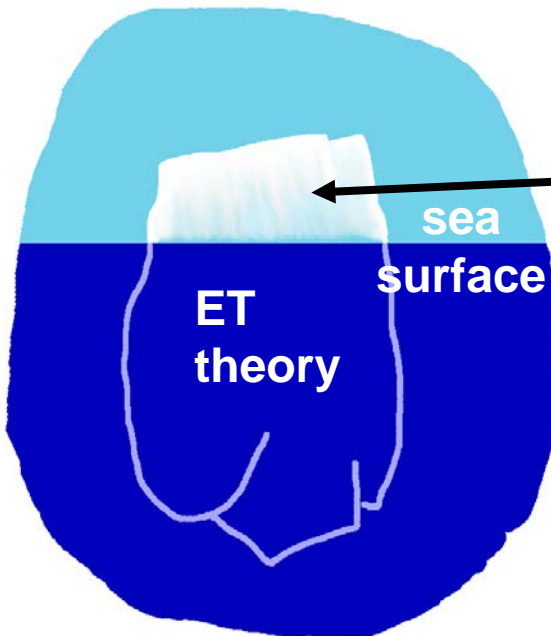
$$\epsilon_0 = 78; \quad \epsilon_\infty = 1.8$$

$$\text{Pekar factor} = \left( \frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} \right) = 0.55$$

$$E_r = \frac{1}{2} \left( \frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} \right) \left( \frac{1}{a} - \frac{2}{R_{12}} \right)$$

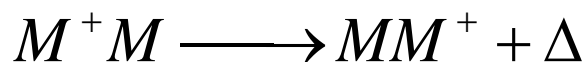
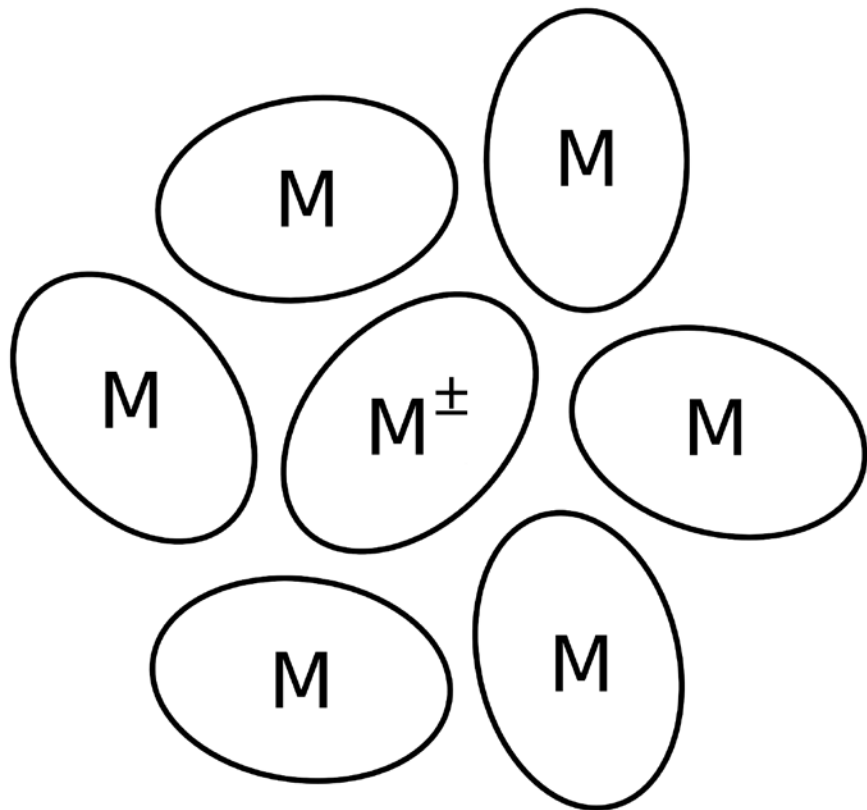
reorganization energy

Levich-  
-Dogonadze  
(1959)  
(Theory of  
multi-phonon  
transitions)



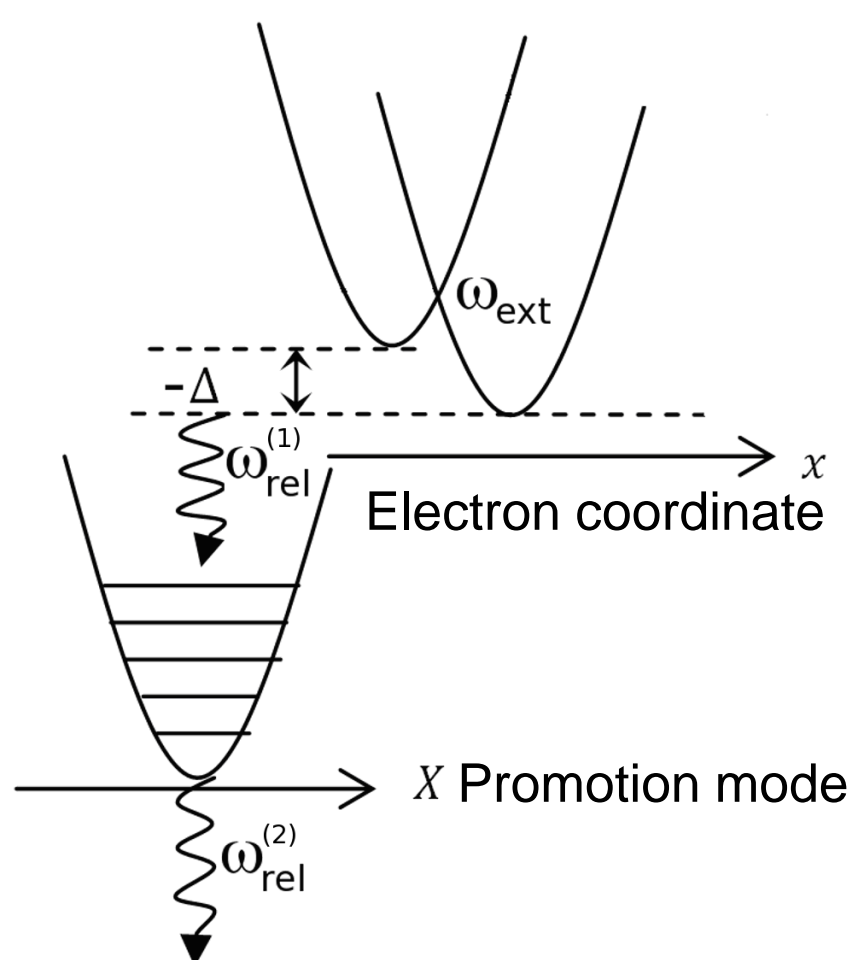
Marcus (1956)  
(phenomemnological)

The ET scheme. The reaction center



$$K(T) \ll \begin{pmatrix} \omega_{ext} \\ \omega_{rel} \end{pmatrix} \ll \omega_{int} \approx \omega_0$$

The energy exchange by means of the promotion mode

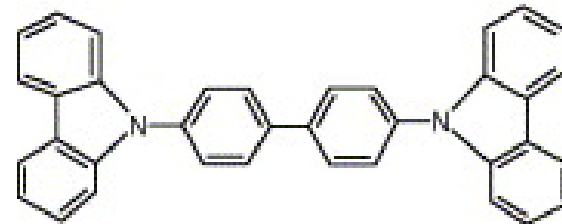
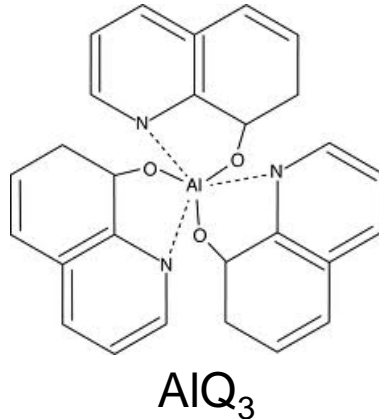


# 3. Charge carrier mobilities in OLED materials

OLED = Organic Light Emitting Diode

Active ET centers of OLED systems are dimers  $(M^\pm)M$  appearing in applied electric fields

Typical monomer (M) molecules



Dielectric permittivity for  $AlQ_3$ :

Static:  $\epsilon_s = 2.84$  (calculated)

or  $3 \pm 0.3$  (experiment)\*)

Optical:  $2 < \epsilon < 3$ ; Pekar factor  $< 0.1$

Conclusion: the ET in OLED active centers is mainly associated with local molecular modes, rather than with medium polarization modes, as in usual ET theories

\*) V. Ruhle et al, *JCTC* **7** (2011) 3335-3345

# 4. The ET without medium polarization

$x$  – essentially quantum coordinate (two states 1 and 2)

$X(X_k)$  – local molecular modes (n or n' states: 1n and 2n')

$Q(Q_\nu)$  – medium modes with continuous frequency spectrum (the bath)

Both polarization and  
acoustic phonon modes may  
be included in Q

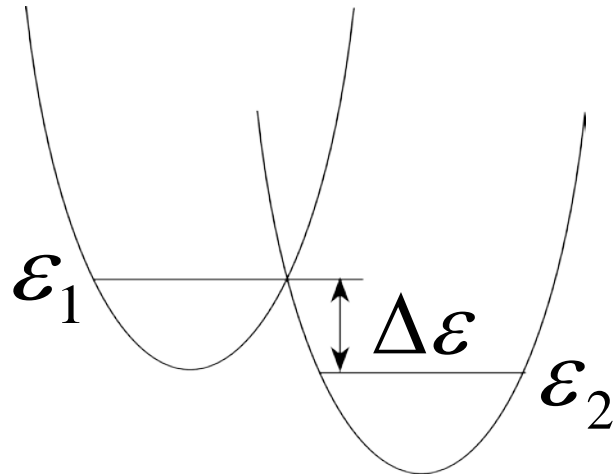
MLD – Marcus-Levich-Dogonadze (traditional ET)

DJ – Dogonadze-Jortner (including local modes)

The model	The interaction scheme	Comment
MLD (spin-boson)	$\begin{array}{c} x \quad X \\ \quad \backslash \\ \quad \quad Q \end{array}$	Strong x-Q, Q - polarization
DJ (generalized spin-boson)	$\begin{array}{c} x \text{---} X \\ \quad \backslash \\ \quad \quad Q \end{array}$	Strong x-Q and x-X, Q – polarization. The same bath Q for all transitions 1n→2n'
Non-spin-boson (the present work)	$\begin{array}{c} x \text{---} X \\ \quad \quad \backslash \\ \quad \quad \quad Q \end{array}$	Strong x-X and weak X-Q, Q – acoustic phonons, X – dependent bath (XQ)

***The connection to the continuum bath Q is obligatory in order to dissipate the energy misfit  $\Delta$  of a ET reaction. This assures the convergence of rate integrals.***

# 5. Fermi Golden Rate



$P_{12}$  (transition probability per time unit)

$$P_{12} = \frac{2\pi}{\hbar} |V_{12}|^2 \delta(\varepsilon_1 - \varepsilon_2)$$

(purely incoherent transitions)

In the medium with continuum frequency spectrum

$$P_{12} = \frac{2\pi}{\hbar} |V_{12}|^2 \rho(\Delta\varepsilon)$$

distribution of the energy misfit

Coherent (resonance) transitions are neglected

Spin-boson model: two levels ( $\varepsilon_1$  and  $\varepsilon_2$ ) in the continuum medium bath, which accepts the energy misfit  $\Delta\varepsilon$

# 6. Active ET local motions:

## Reorganization mode X

(intramolecular) with frequency  $\omega_0$

Transfer

integral:

$$J_{m'} = A_0 \langle \varphi_n(X) | \hat{J}_X | \varphi_n(X') \rangle \quad \hat{J}_X = \exp\left(-\delta \frac{d}{dX}\right) \quad (\text{shift operator})$$

$\varphi_n(X)$  are oscillator functions:

$$J_{m'} = A_0 \langle \varphi_n(X) | \varphi_n(X + \delta) \rangle$$

$$E_r = \frac{m\omega_0^2}{2} \delta^2$$

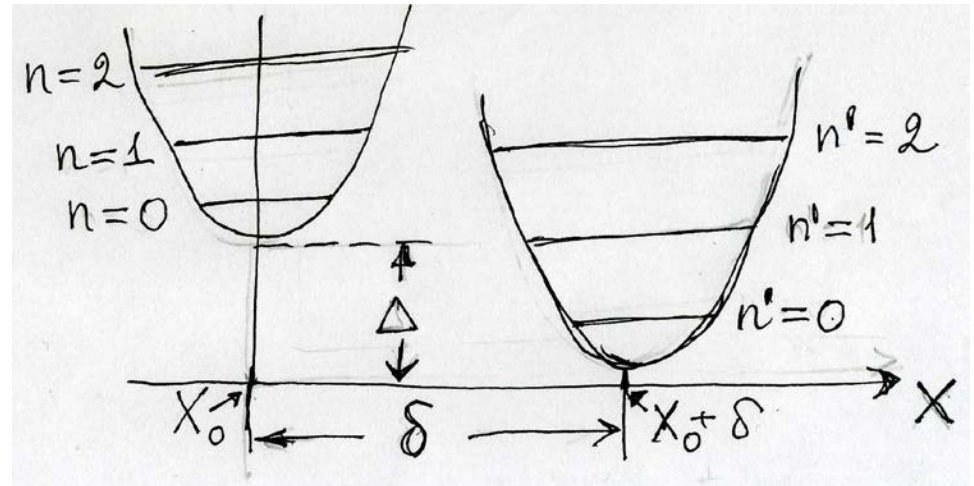
(reorganization energy)

$A_0$  and  $E_r$  (or  $\delta$ ) are  
the basic parameters

Marcus (1956; M)

Levich, Dogonadze (1959; LD)

MLD mechanism of ET



$\delta$  is the shift of the equilibrium position  $X_0$

$\Delta$  is the reaction energy change

# 7. Active ET local motions: promotion

mode  $X$  (intermolecular) with frequency  $\omega_0$

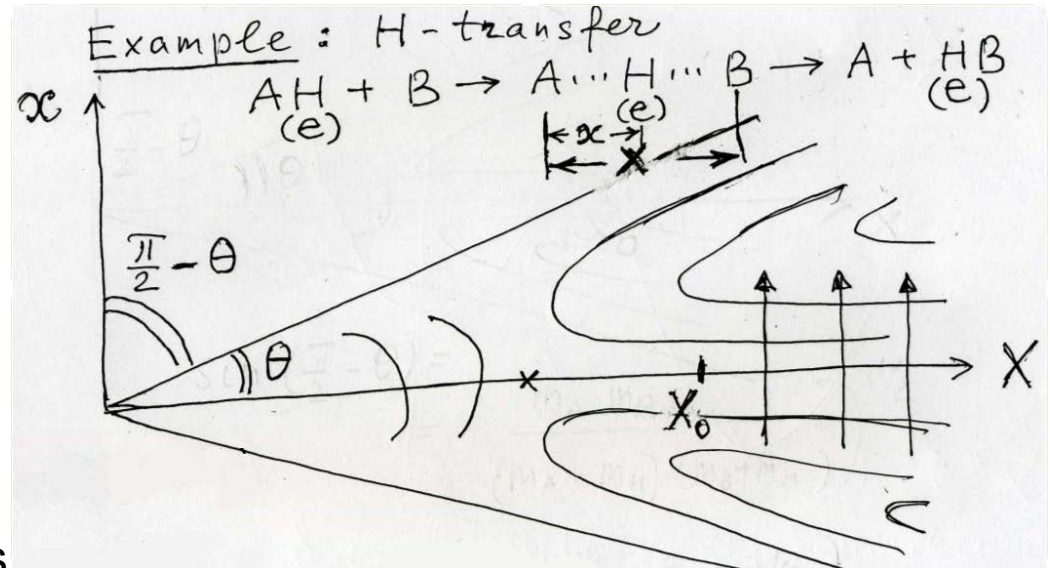
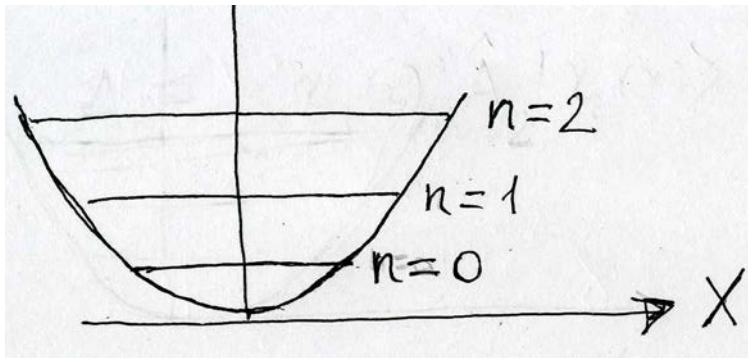
Transfer integral:

$$J_{mn'} = \langle \chi_1(x) \varphi_n(X) | \hat{J} | \chi_2(x) \varphi_{n'}(X) \rangle = A_0 \langle \varphi_n(X) | \hat{J}_X | \varphi_{n'}(X) \rangle$$

$\chi_i(X)$  are electron (or H) functions

$$\hat{J}_X = \exp(-\mu(X - X_0) - \nu(X - X_0)^2)$$

$\varphi_n(X)$  are oscillator functions



$A_0$ ,  $\mu$  and  $\nu$  are the basic parameters

Miller, Abrahams (1960; MA  $\rightarrow$  ET)

Trakhtenberg, Klochikhin, Pshezhetski (1982; T  $\rightarrow$  H transfer)

MAT mechanism of ET or HT

$$\sin\left(\frac{\pi}{2} - \theta\right) = \left[ \frac{m_A m_B}{(m_A + m_H)(m_B + m_H)} \right]^{1/2}$$

$$\theta = 8^\circ \quad (\text{H transfer})$$

$$\theta = 0.6^\circ \quad (ET)$$



# 8. The ET kinetics (Levich, Dogonadze, 1959)

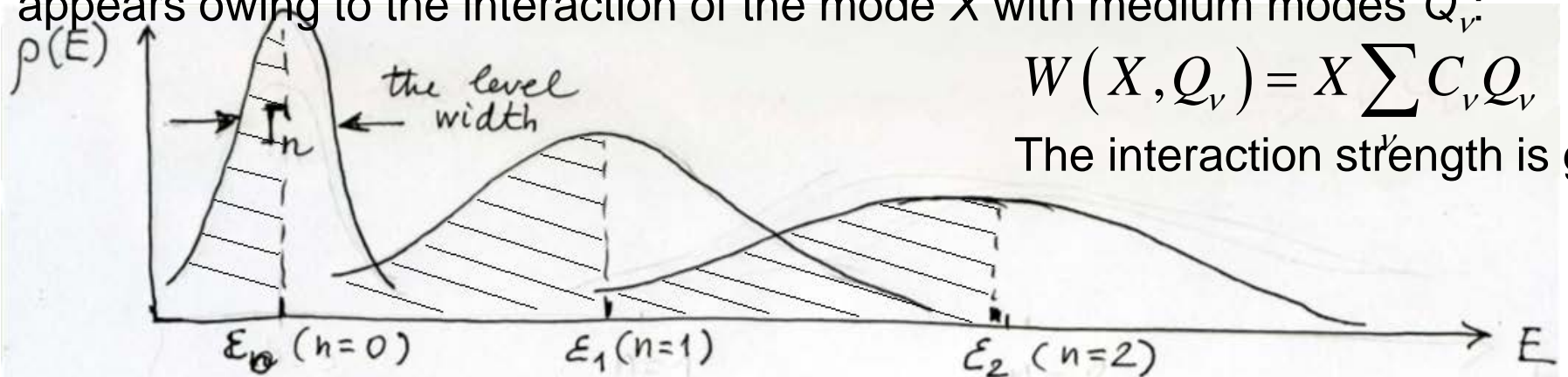
$$K(T) = \frac{2\pi C(T)}{\hbar Z(T)} \quad (1) \quad \text{Eq. (1) is derived from the Fermi "Golden Rule"}$$

The present work follows our approach suggested in 2006:

$$C(T) = \sum_{n,n'} \int_0^\infty dE \exp\left(-\frac{E}{k_B T}\right) |J_{nn'}^2| \rho_n(E) \rho_{n'}(E) \rightarrow \text{the reaction probability flux}$$

$$Z(T) = \sum_n \int_0^\infty dE \exp\left(-\frac{E}{k_B T}\right) \rho_n(E) \rightarrow \text{the partition function}$$

The energy distributions  $\rho_n(E)$  are the basic quantities. The level broadening appears owing to the interaction of the mode  $X$  with medium modes  $Q_v$ :



$$W(X, Q_v) = X \sum C_v Q_v$$

The interaction strength is  $g$

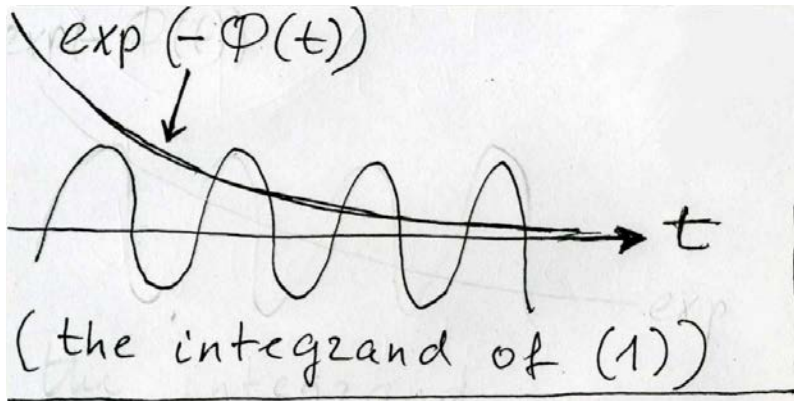
$$\epsilon_n = \left(n + \frac{1}{2}\right) \hbar \omega_0; \quad \Gamma_n = \left[ (2n+1) \coth \frac{\hbar \omega_0}{k_B T} - 1 \right] g$$

*n-dependent continuous frequency spectrum destroys the spin-boson model*

# 9. ET dynamics/kinetics

$$K(T) = \text{const} \int_{-\infty}^{+\infty} \cos\left(\frac{\Delta}{\hbar\omega_0}t\right) \exp(-\Phi(t)) dt \quad (1)$$

Rate constant is determined by the dissipation of the energy misfit  $\Delta$   
 $\omega_0$  is the frequency;  $\Delta$  is the reaction energy change



Invoking the promotion mode  $X$  is quite unusual in the ET theory (i.e. the MA mechanism is usually disregarded)

$\Phi(t)$  is extremely complicated in the full theory (LD, 1959; the reorganization mode  $X$  with frequency  $\omega_0$ )

Marcus (1956; the reorganization mode  $X$ ):

$$K(T) = \frac{A_0^2}{\hbar} \sqrt{\frac{\pi}{E_r k_B T}} \times \exp\left[-\frac{(\Delta + E_r)^2}{4k_B T E_r}\right] \quad (2)$$

Eq. (2) is the asymptotic limit of (1), purely classical, i.e.  $\hbar \omega_0/kT \ll 1$

# 10. The spectral functions $\rho_n(t)$ and $f(t)$

$$\rho_n(E) = \rho_n(\omega) = \int_{-\infty}^{+\infty} \rho_n(t) \exp(i\omega t) dt; \quad \omega = \frac{E - \varepsilon_n}{\hbar\omega_0}$$

(the energy level distribution)

$$\rho_n(t) = \frac{1}{2\pi} \exp\left[-\frac{\Gamma_n}{2\hbar\omega_0} f(t)\right] \quad (\text{the spectral function}) \quad (1)$$

$$f(t) = |t| + \frac{1}{b} (e^{-b|t|} - 1) \quad (\text{Kubo, Toda, Hashitsume, 1978, 1986})$$

$$\Gamma_n = \left[ (2n+1) \coth \frac{\hbar\omega_0}{2k_B T} - 1 \right] g \quad \tilde{\varepsilon}_n = \left( n + \frac{1}{2} \right) \hbar\omega_0 - i \frac{\Gamma_n}{2} \quad (2)$$

Eq. (2) is derived based on the quantum relaxation equation (Bloch,

Redfield);  $g$  is the strength of  $X$ /medium interaction:

$$W(X, Q_\nu) = X \sum_\nu C_\nu Q_\nu = XQ, \quad Q = \sum_\nu C_\nu Q_\nu \quad (3)$$

*Parameters  $b$  (Eq. (1)) and  $g$  (Eq. (2)) can be extracted from the correlation function  $C(t) = \langle Q(t=0)^* Q(t) \rangle$ , where  $Q(t)$  Eq. (3)) is the collective medium coordinate (the medium induced random force).*

# 11. The basic parameters

$\omega_0$  – the frequency of the local mode  $X$

$J_0$  – the transfer integral

$E_r$  – the reorganization energy  $\frac{E_r}{\hbar\omega_0}$

$\Delta = E_2 - E_1$  – the reaction energy change  $\left(\frac{\Delta}{\hbar\omega_0}\right)$  } for the reorganization mode  $X$

$\mu, \nu$  – the parameters of the transfer

integral  $J = J_0 \exp(-\mu X - \nu X^2)$

$\Delta = E_2 - E_1$  – the reaction energy change  $\left(\frac{\Delta}{\hbar\omega_0}\right)$  } for the promotion mode  $X$

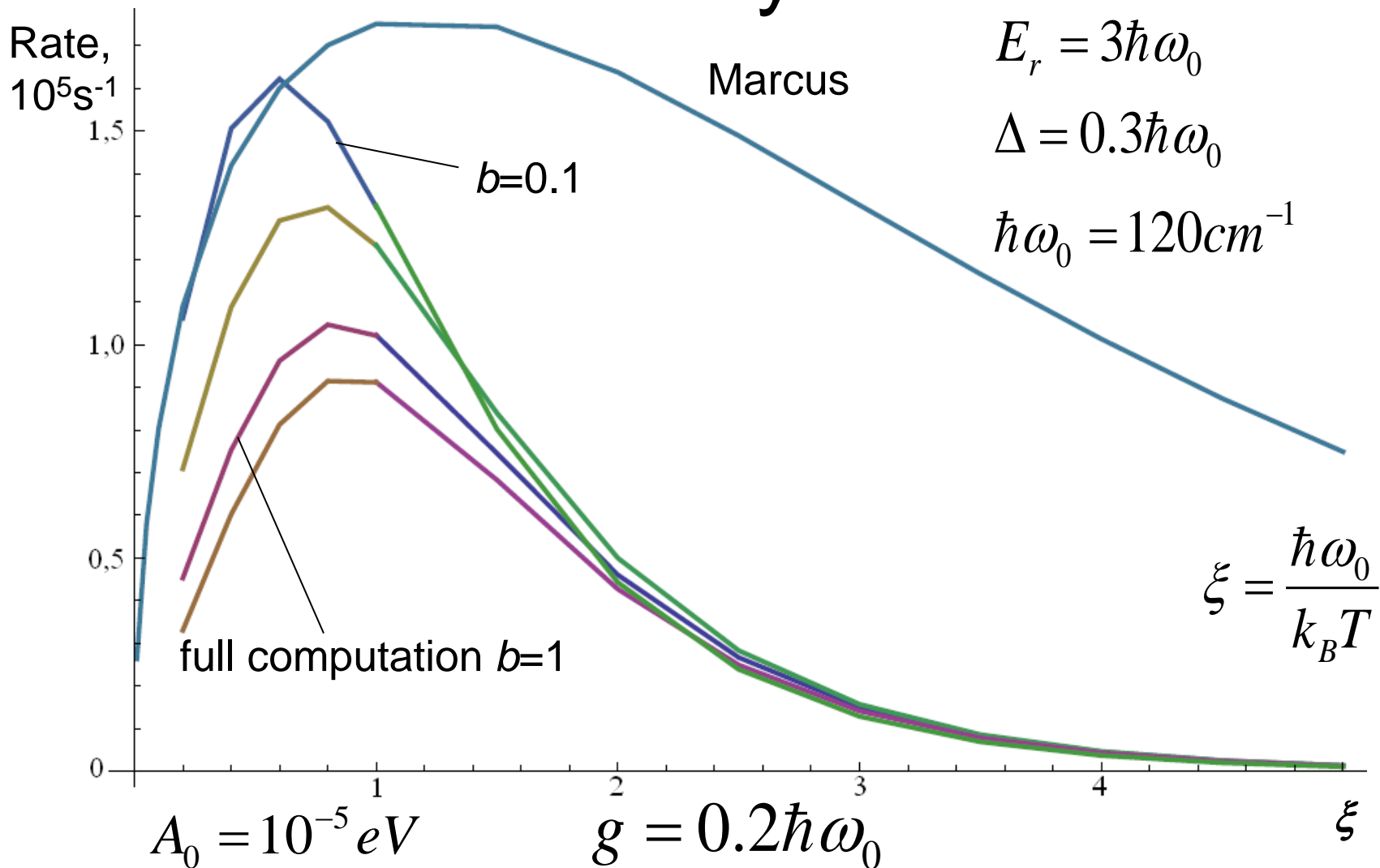
$g$  – the strength of the mode/medium interaction

$b$  – the parameter of Kubo function  $f(t) = |t| + \frac{1}{b}(\exp(-b|t|) - 1)$  } specific for the present approach

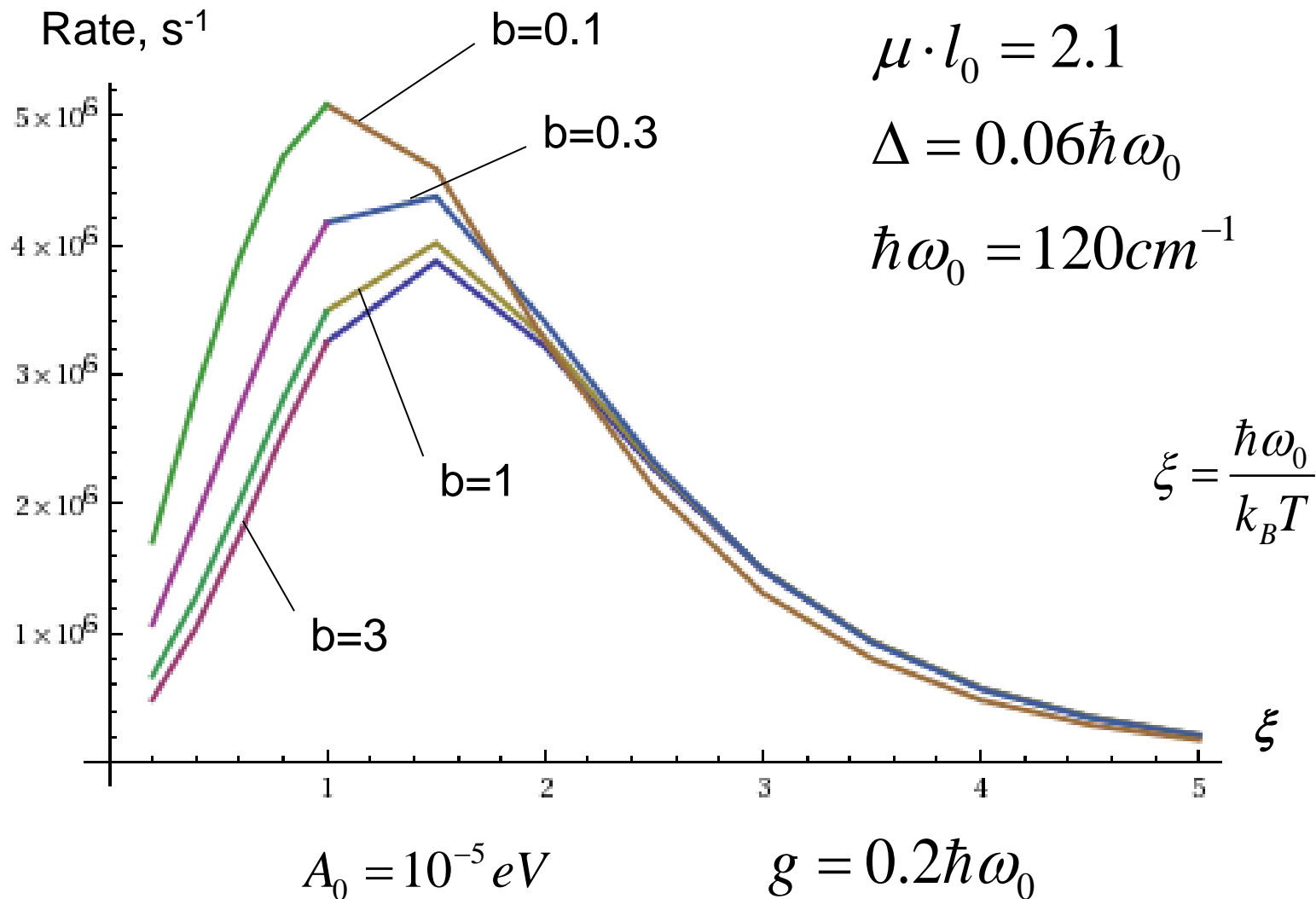
The important parameter  $\xi = \frac{\hbar\omega_0}{k_B T}$

determines the kinetic regime in the whole temperature range

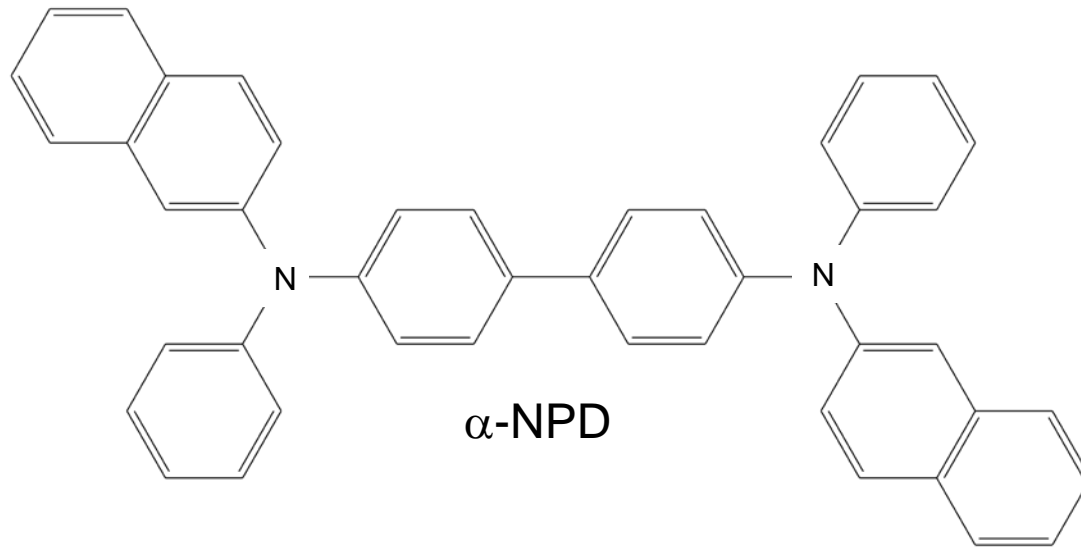
# 12. The ET kinetics in typical OLED systems



# 13. The ET kinetics for promotion mode

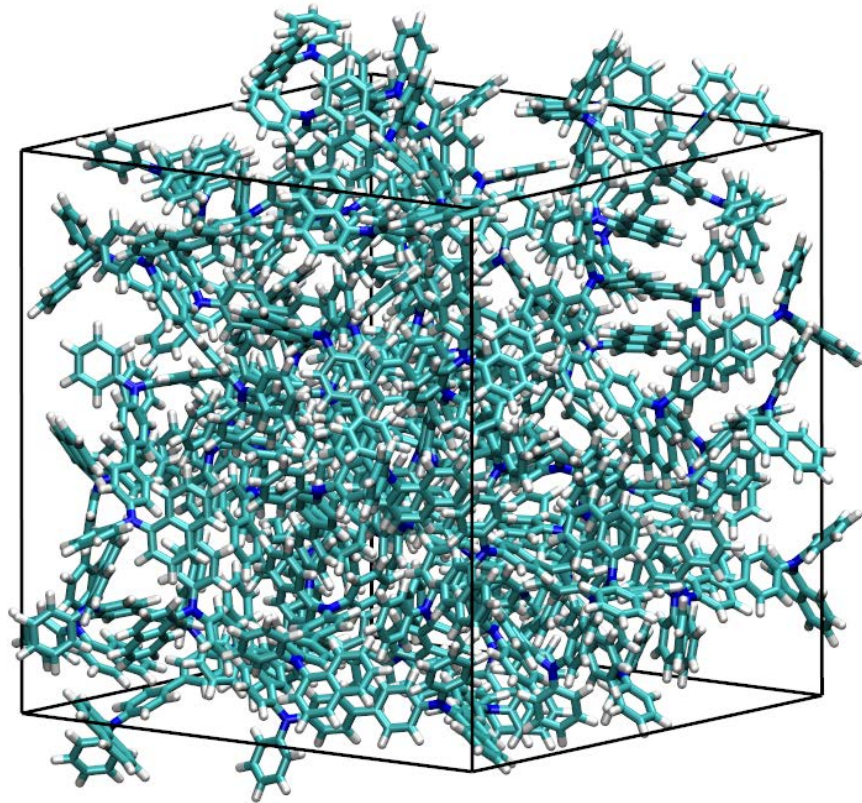


# 14. The conduction material



N,N – di[1-heptaphenyl – N,N-diphenyl] – 1,1 biphenyl – 4,4 diamine

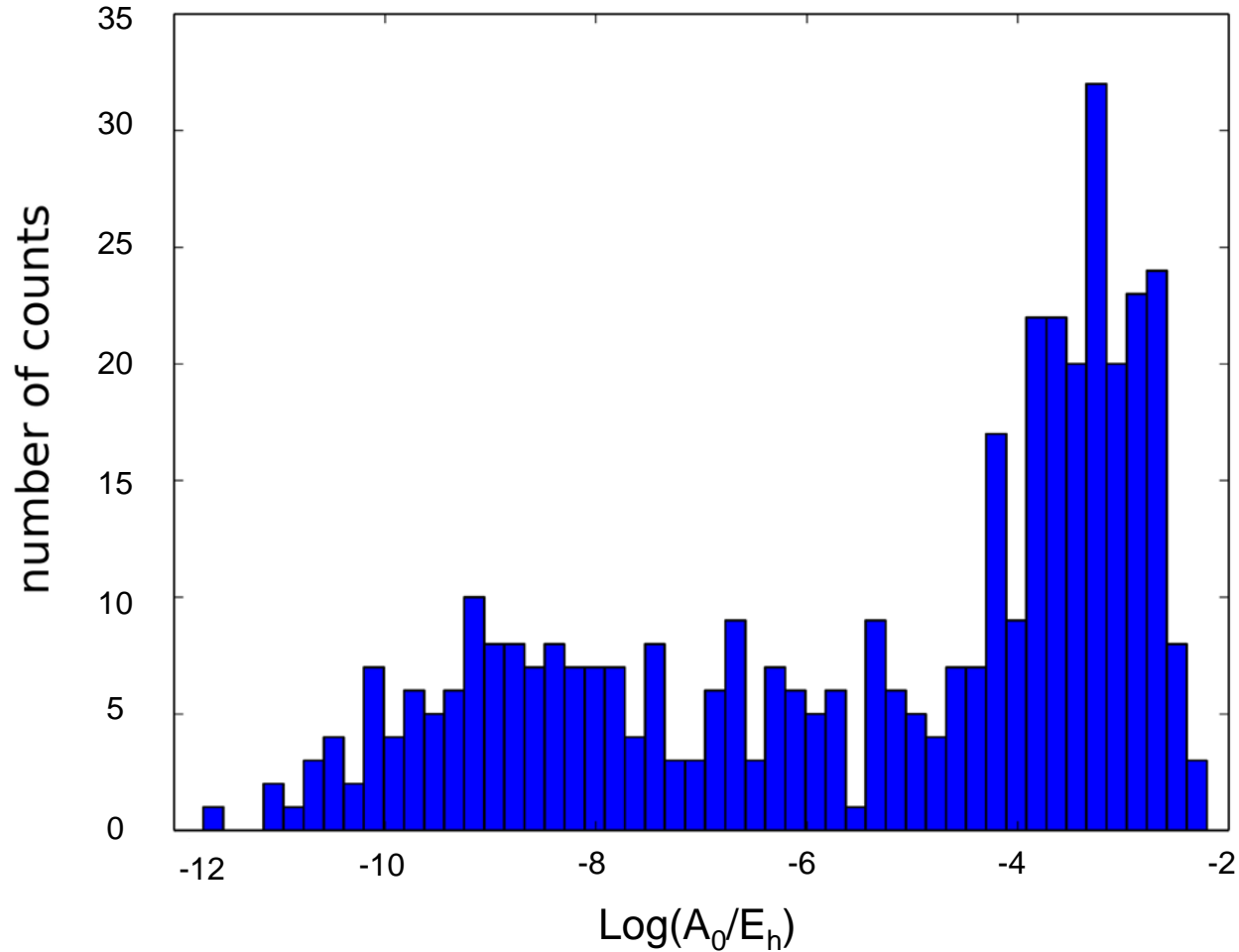
# 15. The MD simulation of amorphous $\alpha$ -NPD material.



the MD cell contains 50  $\alpha$ -NPD molecules



# 16. The distribution of transfer integrals $A_0$



# 17. The computations for eight $\alpha$ -NPD dimers

$N_0$	$A_0$ , $10^{-2}$ eV	$\mu$ , $\text{nm}^{-1}$	$\omega_0$ , $\text{cm}^{-1}$	$b/\omega_0$	$g/\omega_0$	$\langle x^2 \rangle$ , $\text{nm}^2$	$\Delta$ , eV	$\log(K_{ij})$ , $\text{s}^{-1}$
1	0.85	20.1	155	47.8	1.12	0.021	-0.53	12.66
2	1.51	17.0	218	12.2	0.86	0.020	-0.22	13.33
3	1.09	16.0	183	17.6	1.04	0.023	-0.08	12.75
4	2.68	28.1	190	37.6	0.75	0.019	-0.40	15.51
5	2.29	3.4	187	12.0	1.01	0.019	-0.76	12.21
6	2.55	13.6	236	11.2	0.80	0.019	-0.24	13.34
7	0.93	23.2	149	22.3	1.15	0.034	-0.18	14.24
8	1.94	34.6	158	20.8	0.92	0.022	-1.24	16.50

$A_0$ : transfer integrals (QC)     $\Delta$ : the energy misfit (QC)     $K_{ij}$ : the calculated rate constant

$$\mu: A_0(X) = A_0 \exp[-\mu(X - X_0)]$$

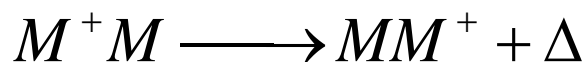
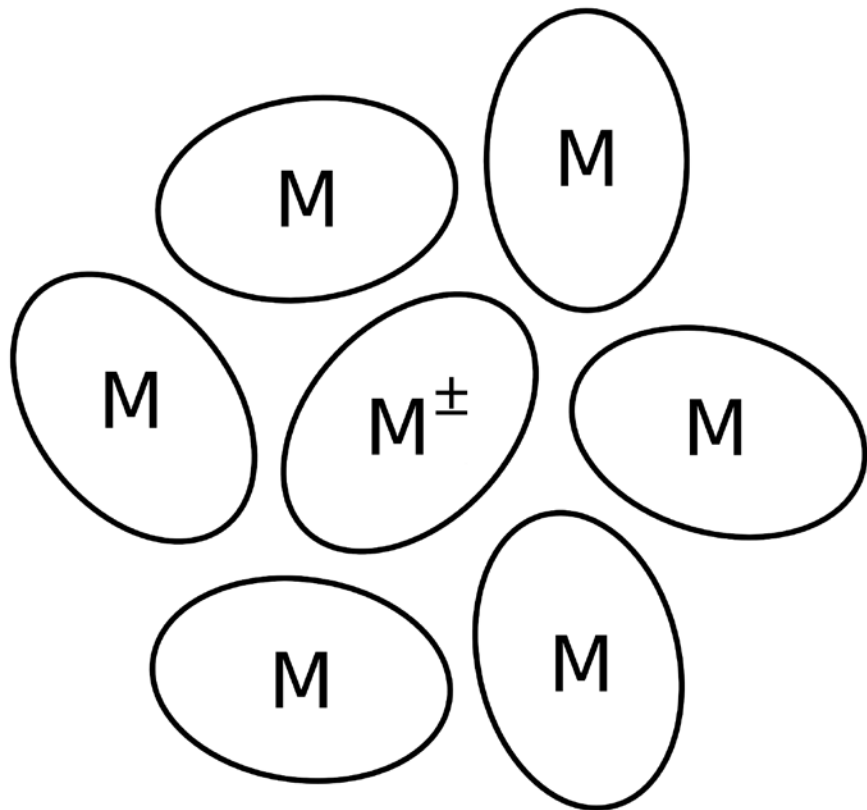
$\omega_0$ : the frequency  
of promotion mode

$b, g$ : the relaxation  
parameters

MD simulation of the velocity

correlation function     $\langle \dot{X}(t=0) \cdot \dot{X}(t) \rangle$

The ET scheme. The reaction center



$$K(T) \ll \begin{pmatrix} \omega_{ext} \\ \omega_{rel} \end{pmatrix} \ll \omega_{int} \approx \omega_0$$

The energy exchange by means of the promotion mode

